

## Inertial propeller (Theory)

The offered material is continuation of clause of the author: “The Third law of Newton is not carried out for a unbalanced body with rotary fluctuation” // - “Гравитон” # 12, 2005, p. 9.

The principle of operation of the device “Inertial propeller” can be explained using the scheme without a gyroscope (see Fig 1).

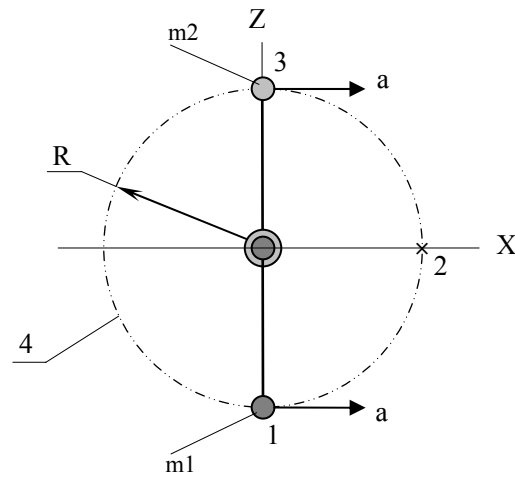


Fig 1

Along the ring 4 with radius  $R$ , masses  $m_1$  and  $m_2$  can oppositely-symmetrically move using the power drive. In order to simplify the numerical evaluations, we set  $m_1 = m_2$ , while the sum of masses  $m_1 + m_2 = m$ , where  $m$  is the mass of ring 4 along with the drive. The sequence of movement is as follows. Masses  $m_1$  and  $m_2$  speed up continuously with the tangential acceleration  $a$  within the sectors 1-2 and 3 - 2 respectively. Masses  $m_1$  and  $m_2$  slow down continuously with the tangential acceleration  $-a$  within the sectors 2-3 and 2-1 respectively to a full stop at points 3 and 1. Then, the acceleration/deceleration process is repeated in the opposite direction. The point 2 is the acceleration reversal one. In order that the acceleration/deceleration process started and was over at points 1 and 3, the braking moment  $M_T$  should be taken to be equal to the acceleration moment:  $M_T = Ma$ . For definiteness sake, let the braking torque  $M_T$  be the frictional moment.

So, under the action of the acceleration moment  $Ma$ , the acceleration of masses  $m_1$  and  $m_2$  begins at initial points 1 and 3 respectively and ends at point 2 where the rotary drive is turned off ( $Ma = 0$ ). At that instant, the pulse  $P$  of mass  $m$  is equal to zero:  $P = 0$ , and each of masses  $m_1$  and  $m_2$  takes on the tangential pulse  $P_m \neq 0$ .

Further, let us assume first that the braking moment is absent within the sectors 2-3 and 2-1 ( $M_T = 0$ ).

The subsequent movement of masses  $m_1$  and  $m_2$  occurs mechanically. At that, under the action of the centripetal force and in conformity with the law of action and reaction, the turning of pulse vectors  $P_m$  will occur opposite to the  $X$  axis direction (without change in their values). At the same time, the vector of pulse  $P$

will increase along the X axis direction. When masses  $m_1$  and  $m_2$  move to the points 3 and 1, these pulses will become numerically equal and opposite in direction ( $2P_m = P$ ). In this case, the system common centre of gravity will remain motionless.

Let us repeat the interaction process beginning from the point 2 with turned-on braking moment,  $M_T \neq 0$  (at the same time,  $M_a = 0$ ). For definiteness sake, the braking moment  $M_T$  is the moment of dry friction numerically equal to

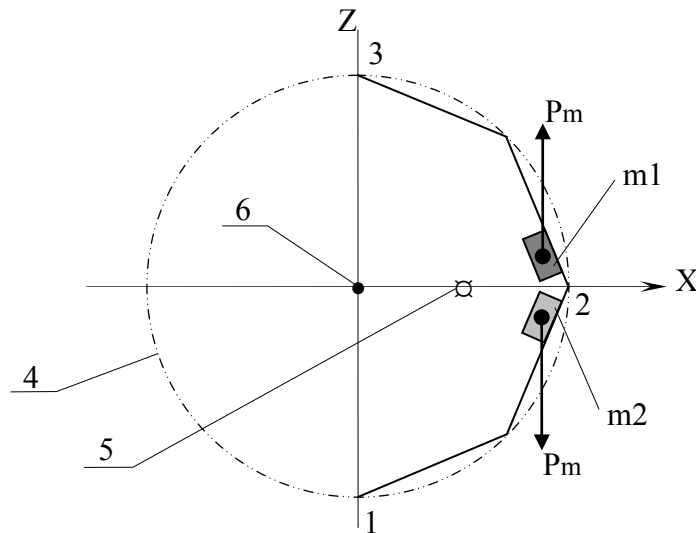


Fig 2

the accelerating moment:  $M_T = M_a$ .

So, prior to beginning of slowing down of masses  $m_1$  and  $m_2$ , each of them has taken on the tangential pulse equal to  $P_m$ . In this case, the pulse  $P$  of the rest of mass  $m$  of the system is equal to zero:  $P = 0$  as mass  $m$  was motionless at this instant. The subsequent movement of masses  $m_1$  and  $m_2$  occurs mechanically

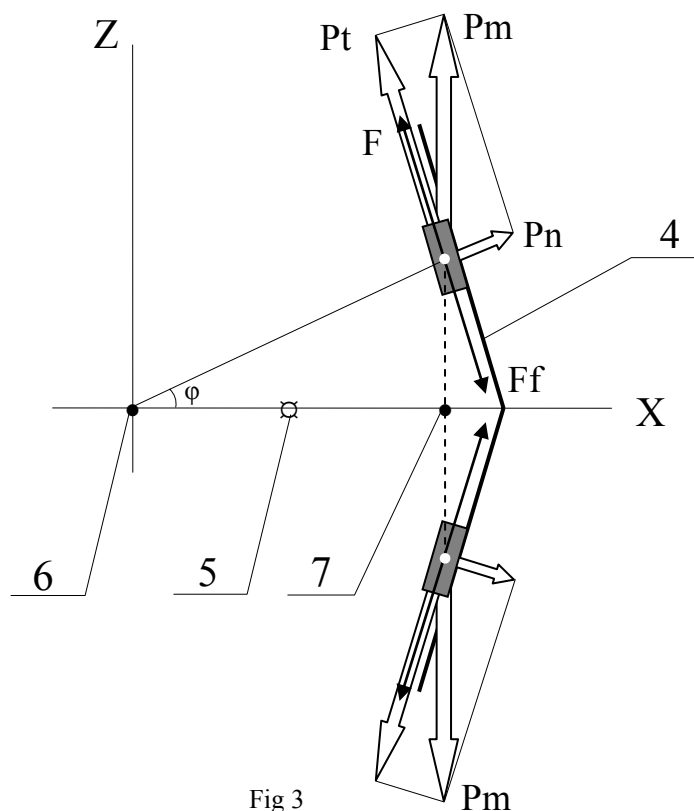


Fig 3

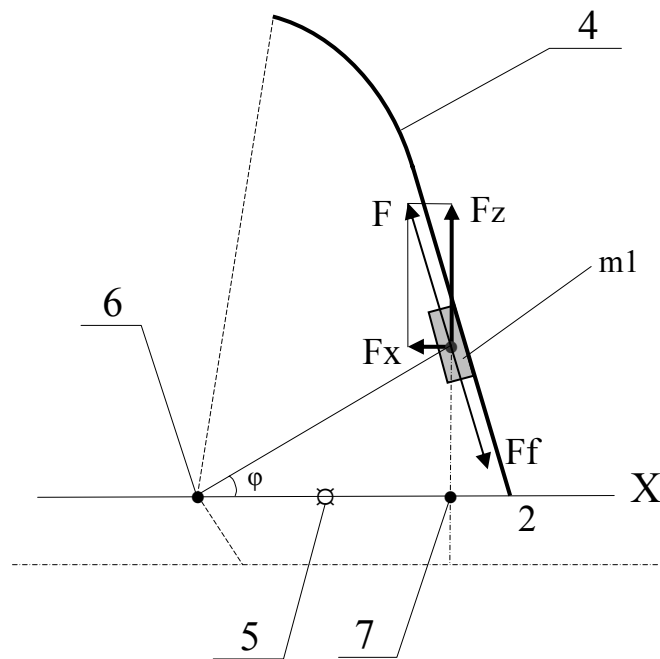


Fig 4

but their motion is hindered by the friction force  $F_f$  on the mass  $m$  side which generates the frictional moment  $M_T$  (it is the braking moment). For the greater obviousness sake, we represent a part of the ring 4 in the form of inclined planes as shown in Fig.2 and enlarged in Fig 3. Hereinafter: 6 - center of gravity for mass  $m$ , 7 - center of gravity for sum  $(m_1 + m_2)$ , 5 - common center of gravity for the whole system. The vector of pulse  $P_m$  can be resolved into two components: parallel to the plane -  $P_t$  and perpendicular to the plane -  $P_n$ . In the course of sliding of masses  $m_1$  and  $m_2$  over the planes 4 against the friction forces  $F_f$ , the work is done, therefore, the magnitude of pulse  $P_t$  will decrease. However, as it is well known, the change in the magnitude of pulse is the force  $F = \partial(P_t)/\partial t$  where  $\partial(P_t)/\partial t$  is the derivative of pulse  $P_t$  with respect to time  $t$ . The force  $F$  is numerically equal to the friction force:  $F = F_f$  and is in opposition to the latter. The force  $F$  is force with which the masses  $(m_1 + m_2)$  act on mass  $m$ . By its nature, it is inertial force. Let us resolve the force  $F$  for each mass  $m_1$  and  $m_2$  into two components:  $F_x$  and  $F_z$ . The forces  $F_z$  are mutually balanced while the forces  $F_x$  are not balanced. Masses  $m_1$  and  $m_2$  together with the ring 4 form the common mechanical system, therefore, the forces  $F_x$  are applied to the common center of gravity 5 of the system. Hence it follows that the center of gravity 5 of the whole system is shifted in the direction of action of forces  $F_x$ . The described interaction is similar to that of two bodies in case of the oblique inelastic impact with each other when, initially, one of bodies was motionless ( $m$ ) while the other ( $m_1 + m_2$ ) has moved (i.e., had a pulse). After collision, these bodies stick together and continue the motion as a single whole. Fig 5 and Fig 6 present the illustrative example of interaction of two plasticine balls 8 imitating the direct impact as a result of which the common center of gravity 5 is shifted. Position 9 in Fig 6 shows the location of the common center of gravity 5 of the

system prior to interaction.

The unit “Inertial Propeller” (Pub. No. US 2005/0169756 A1) differs from one described in the text in that one of masses, for example,  $m_2$ , is replaced by a gyroscope. In the unit “Inertial Propeller” (see Pub. No. US 2005/0169756 A1, Fig.3), the points  $a$ ,  $b$  and  $c$  correspond to points 1, 2 and 3 of the current description. The mass of imbalance 7 corresponds to mass  $m_1$ . The starting moment  $M_\Pi$  corresponds to the accelerating moment  $M_a$  while the counter moment ( $-M_\Pi$ ) corresponds to the braking moment  $M_T$  in the current description. At the point  $b$  of the unit “Inertial Propeller”, there is a reversal of the rotary moment of imbalance 7: the accelerating moment  $M_\Pi$  is turned off and the braking moment ( $-M_\Pi$ ) is turned on. It is evident that the braking moment ( $-M_\Pi$ )  $\equiv$   $M_T$  can be of any nature: it can be moment of dry friction as shown above or electrodynamic moment originating as a result of the electric motor operation, for example, in the mode of current generation.

It should be noted once more (see Pub. No. US 2005/0169756 A1, Fig 3) that, under the action of the drive moment  $M_\Pi$ , the acceleration of imbalance 7 occurs within the sector  $a - b$ . At point  $b$ , the moment  $M_\Pi$  is turned off (i.e.  $M_\Pi = 0$ ). The position of imbalance 7 at point  $b$  corresponds to the initial position of plasticine balls 8 in Fig. 5 of the current description. At that, imbalance 7 takes on the tangential pulse  $P_m = m_1 \cdot \mathbf{a} \cdot t = m_1 \sqrt{2\varphi R \mathbf{a}}$ , where  $m_1$  is the mass of imbalance 7,  $\mathbf{a}$  is tangential acceleration of imbalance,  $t$  is a duration of turning from point  $a$  to point  $b$  and  $\varphi$  is angle to the X axis,  $R$  is the radius of rotation. The pulse  $P$  of the rest of mass ( $m$ ) of the unit at this instant was equal to zero as the mass  $m$  at that moment is motionless (position 6 in Fig 5 hereof). At that, the numerical value of tangential pulse of imbalance 7 at point  $b$  is  $P_m = m_1 \sqrt{\pi R \mathbf{a}}$ . The subsequent motion of imbalance 7 in the sector  $b - c$  occurs mechanically but, at that, the centripetal force on the mass  $m$  side and braking moment ( $-M_\Pi$ ) on the side of the electric motor of rotation drive are applied to it. As a result of turning of the vector  $P_m$  according to the law of action and reaction, its projection onto X axis increases:

$$P_x = P_m \cdot \sin\varphi = m_1 \sqrt{\pi R \mathbf{a}} \cdot \sin\varphi,$$

and, at the same time, the pulse  $P$  of the rest of mass ( $m$ ) in the opposite direction increases which is equal to:

$$P = m_1 \sqrt{2\varphi R \mathbf{a}} \cdot \sin\varphi.$$

In this case, the centers of gravities 6 and 7 draw together while the common center of gravity 5 remains motionless (see Fig 5 hereof). Within the range of angles of  $\pi/2 > \varphi > 0$ ,  $P_x > P$ , i.e.:

$$\Delta P = P_x - P = m_1 \sqrt{\pi R \mathbf{a}} \cdot \sin\varphi - m_1 \sqrt{2\varphi R \mathbf{a}} \cdot \sin\varphi \neq 0.$$

Because, at that, the value of the pulse  $P_m$  decreases simultaneously under the action of the braking moment ( $-M_\Pi$ ) then, along the X axis, the component of force  $F$  appears,  $F_x = \partial(P_x)/\partial t$  (where  $P_x$  is the projection of vector  $P_t$  onto the X axis, see Fig 3 hereof), which is applied to the common center of gravity 5 of the whole mechanical system.

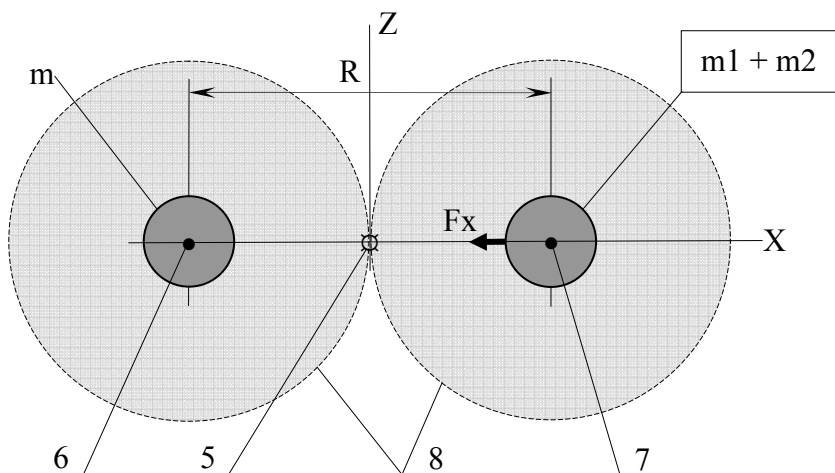


Fig 5

The force  $F_x$  shifts the common center of gravity 5 of the system in the same manner as shown in Fig 5 - Fig 6 hereof and well known from experiments of inelastic collision of, for example, plasticine balls 8.

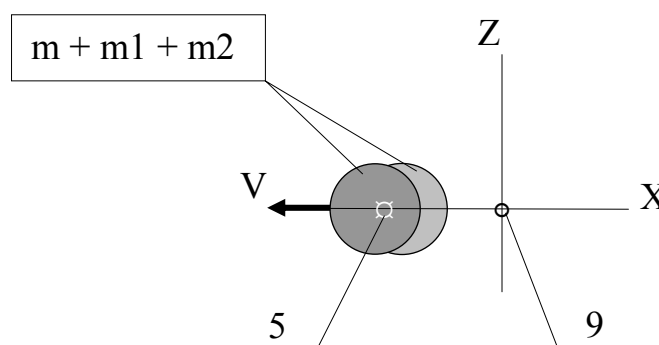


Fig 6

## THE LITERATURE

1. Линеви́ч Э. И. «Геометрическое обоснование эксперимента Хаясака-Такеучи с вращающимися роторами». Доклад на 2-ой СНГ Межнаучной конференции «Единая теория мира и её практическое применение». 20-21 сентября 1993 г. Петрозаводск, Россия.
2. Hayasaka H., Takeuchi S. Phys. Rev. Lett.– V.63. P.2701-2704.
3. Линеви́ч Э. И. Явление антигравитации физических тел (ЯАФТ).– Хабаровск: ПКП «Март», 1991.
4. Линеви́ч Э. И. Динамическая симметрия вселенной.– «Природа и аномальные явления» № 1-2, 1995, с. 6, г. Владивосток.
5. Линеви́ч Э. И. О технической возможности управления темпом времени.– «Гравитон» №8, 2002, с. 10-11.
6. Kishkintsev V. A. Galilean Electrodynamics, 1993. V.4, №3, P.47-50.
7. Forward R. L. Journal of Propulsion and Power. 1989 №1, p.28-37.
8. Линеви́ч Э. И. Аналитический вывод физических констант на основе классических представлений.– Ноябрь 1999 (в переписке с ред. «Гравитон» и с [bradleu@usra.edu](mailto:bradleu@usra.edu)), или <http://www.dlinevitch.narod.ru/analitika.htm>

9. Линеви́ч Э. И., Ежов А. Ф. Инерционный движитель.– «Новая энергетика» №3, 2004, с. 12-15.
10. Линеви́ч Э. И. Гравиинерционный двигатель. Патент RU2080483. 4.05.1994
11. Астахов А. В., Широков Ю. М. Курс физики т.3. Квантовая физика/ Под ред. Ю. М. Широкова.– М.: Наука, 1983.
12. Шипов Г. И. Теория физического вакуума: Теория, эксперименты и технологии. 2-е изд., испр. и доп.– М.: Наука, 1996.
13. Абрамов И. М., Брехман И. И., Лавров Б. П., Плисс Д. А. «Явление синхронизации вращающихся тел (роторов)». Диплом №333. Журнал «Открытия изобретения» №1, 1988.
14. Калашников С. Г. Электричество.– М., 1977, с. 155.
15. [www.dlinevitch.narod.ru](http://www.dlinevitch.narod.ru)
16. <http://www.ntpo.com/physics/studies/28.shtml>
17. <http://www.sciteclibrary.ru/rus/catalog/pages/3964.html>
18. Яблонский А. А. Курс теоретической механики. Ч.2. Динамика.– М., «Высшая школа», 1971.
19. <http://www.tts.lt/~nara/amper/neutron.html>
20. Толчин В. Н. Инерцоид.– Пермь: Пермское книжное издательство. 1977.
21. Round R. V., Rebka G. A., Phys. Rev. Let., 1960, V.4, P.337.
22. Линеви́ч Э. И. «Третий закон Ньютона не выполняется для неуравновешенного тела с вращательным колебанием»// - «Гравитон» №12, 2005, с. 9.
23. Linevich E. I. On basics of potential dynamics.- «New Energy Technologies» #2, 2005, p.44 – 48.
24. <http://vuz.exponenta.ru/PDF/book/t400.pdf>

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